

# Resolving singularities in higher spin gravity

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work in progress with

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# Overview

- Introduction

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- 3D AdS higher spin gravity

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- Remarks on falloff conditions
- Outlook

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- $SL(N, \mathbb{R}) \times SL(N, \mathbb{R})$  CS theory describes spins  $2, 3, \dots, N$
- asymptotic symmetry algebra is extended from Virasoro to  $\mathcal{W}_N$  algebra (Campoleoni, Fredenhagen, Pfenninger, Theisen 2010)

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- Many CFT states with  $0 \leq L_0 \leq c/24$ . Smooth conical defect solutions?

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- Can higher spins resolve singularities in a similar way?

## 3D higher spin gravity

- $SL(N, \mathbb{R}) \times SL(N, \mathbb{R})$  Chern-Simons formulation

$$S = S_{CS}[A] - S_{CS}[\tilde{A}] ,$$
$$S_{CS}[A] = \frac{k}{4\pi} \operatorname{tr} \int_{\mathcal{M}} \left( A \wedge dA + \frac{2}{3} A \wedge A \wedge A \right) \quad k = \frac{l}{4G}$$

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- gauge invariance

$$\delta A = dA + [A, \lambda], \quad \delta \tilde{A} = d\tilde{A} + [\tilde{A}, \tilde{\lambda}]$$

- choose the principal embedding of  $sl(2, \mathbb{R}) \subset sl(N, \mathbb{R})$

$$L_0 = \frac{1}{2} \begin{pmatrix} (N-1) & 0 & \dots & & 0 \\ \vdots & & & & \\ & & (N+1-2i) & & \\ & & & \ddots & \\ 0 & \dots & & & -(N-1) \end{pmatrix},$$

$$L_1 = \begin{pmatrix} 0 & \dots & \dots & \dots \\ -1 & 0 & \dots & \dots \\ & -1 & 0 & \dots \\ & \dots & \dots & -1 & 0 \end{pmatrix}, \quad L_{-1} = \begin{pmatrix} 0 & N-1 & \dots & & 0 \\ & \vdots & & & \\ & & 0 & i(N-i) & \\ & & & \ddots & \ddots \\ 0 & \dots & & & 0 & (N-1) \\ & & & & 0 & 0 \end{pmatrix}$$



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- remaining generators  $W_m^{(s)}$ ,  $m = -(s-1), \dots, s-1$ .

Transform in  $2s-1$  representations,  $s = 3, \dots, N$ :

$$[L_i, L_j] = (i-j)L_{i+j}, \quad [L_i, W_m^{(s)}] = (i(s-1) - m)W_{i+m}^{(s)}$$

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- other embeddings of  $sl(2, \mathbb{R})$  lead to different spin content

(Ammon, Gutperle, Kraus, Perlmutter; Campoleoni, Fredenhagen, Pfenniger 2011)

- generalized vielbein and spin connection

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$$\delta e = d\xi + [\omega, \xi] + [e, \Lambda],$$

$$\delta \omega = d\Lambda + [\omega, \Lambda] + \frac{1}{l^2} [e, \xi]$$

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- metric-like fields (invariant under frame rotations)

$$g_{\mu\nu} = \frac{1}{\text{tr}L_0^2} \text{tr}(e_\mu e_\nu), \quad \phi_{\mu\nu\rho} \sim \text{tr}(e_{(\mu} e_\nu e_{\rho)}), \dots$$

generalized diffeomorphisms act nontrivially on the metric

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- take  $\mathcal{M}$  topologically  $\mathbb{R} \times D_2$ , coordinates  $(t, \rho, \phi)$ ,  
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- imposing boundary conditions + gauge-fixing

$$\begin{aligned} A &= b^{-1} a_+ b dx^+ + b^{-1} db, & b &= e^{(\rho - \rho_0)L_0} \\ \tilde{A} &= b \tilde{a}_- b^{-1} dx^- + b db^{-1} \end{aligned}$$

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- will take  $a_+, \tilde{a}_-$  constant. Trivial holonomy around  $\phi$ -circle imposes

$$e^{2\pi a_+} = 1 \quad (-1) \quad e^{2\pi \tilde{a}_-} = 1 \quad (-1) \quad N \text{ even}$$

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- nontrivial CS holonomy  $\Rightarrow$  gauge fields are not smooth

# General conical defects

- Ansatz for CS connections ( $W_i^{(2)} \equiv L_i$ ,  $c_s = \sqrt{-\frac{\text{tr} L_0^2}{\text{tr} W_{-1}^{(s)} W_1^{(s)}}}$ )

$$a_+ = \sum_{s=2}^N c_s \left( w_1^s W_1^{(s)} + w_{-1}^s W_{-1}^{(s)} \right)$$

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$$\begin{aligned}
 \frac{ds^2}{l^2} &= d\rho^2 - [e^{2\rho} + \Lambda e^{-2\rho} - 2M] dt^2 + [e^{2\rho} + \Lambda e^{-2\rho} + 2M] d\phi^2 \\
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$$M = -\vec{w}_1 \cdot \vec{w}_{-1} \quad \Lambda = (\vec{w}_1 \cdot \vec{w}_1)(\vec{w}_{-1} \cdot \vec{w}_{-1})$$

- conical defect metrics when

$$\Lambda = M^2 \quad \Leftrightarrow \quad \vec{w}_1 = \alpha \vec{w}_{-1}$$

gauge fields are of the form

$$a_+ = \begin{pmatrix} 0 & \alpha v_1 & & & & & \\ -v_1 & 0 & \alpha v_2 & & & & \\ & -v_2 & 0 & & & & \\ & & & \ddots & & & \\ & & & & & 0 & \alpha v_{N-1} & (0) \\ & & & & & -v_{N-1} & 0 & (0) \\ & & & & & (0) & (0) & (0) \end{pmatrix}$$

# Smooth conical defects

- Smooth  $Sl(N, \mathbb{R})$  connection:  $a_+, \tilde{a}_-$  have eigenvalues

$$\lambda_i = (in_1, -in_1, in_2, -in_2, \dots, in_{[N/2]}, -in_{[N/2]}, (0)) \quad n_i \in \mathbb{Z} \text{ or } \mathbb{Z}/2$$



- mass spectrum

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- energy spectrum and deficit angles

$$L_0 = \tilde{L}_0 = \frac{c}{24}(4M + 1) \quad \delta = 2\pi(1 - \sqrt{-4M})$$

- conical defect spectrum up to  $N = 5$

$N$	$(n_i)$	$M$	$\frac{24}{c} L_0$	$\frac{\delta}{2\pi}$
2	(1/2)	-1/4	0	0
3	(1)	-1/4	0	0
4	(3/2, 1/2)	-1/4	0	0
	(3/2, 0)	-9/40	1/10	$1 - \sqrt{9/10}$
	(1, 1)	-1/5	1/5	$1 - \sqrt{4/5}$
	(1, 0)	-1/10	3/5	$1 - \sqrt{2/5}$
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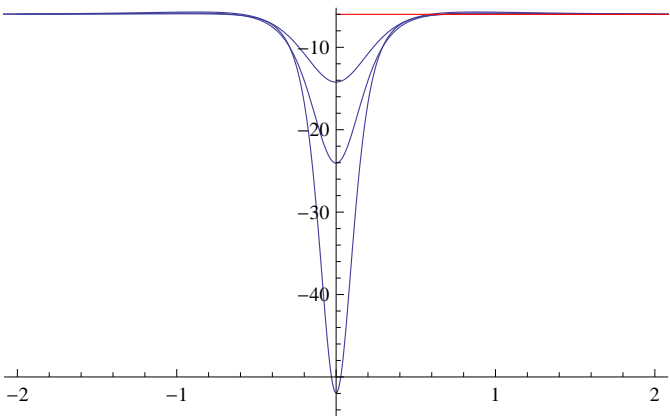
- large  $N$ :  $\mathcal{O}(N^{3N/2})$  conical defects.





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- scalar curvature  
( $N = 5, n_1 = n_2 = \alpha = 1, \gamma_1 = -\gamma_2 = 0, 0.2, 0.3, 0.4$ )



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- Lead to reduced phase space (Drinfeld-Sokolov reduction) with classical  $\mathcal{W}_N$  symmetry
- $\mathcal{W}_N$  survives on quantum level: captures vacuum representation (Gaberdiel, Gopakumar, Saha 2010)

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- Gauge-equivalent?  $u$  can be gauge-fixed such that

$$a_+ = \begin{pmatrix} 0 & u_1 & u_2 & \cdots & u_{N-1} \\ -1 & 0 & 0 & \cdots & 0 \\ & & \cdots & & \\ 0 & 0 & \cdots & -1 & 0 \end{pmatrix}$$

- If eigenvalues not degenerate, diagonalized by Vandermonde

$$V = \begin{pmatrix} \lambda_1^{N-1} & \lambda_2^{N-1} & \cdots & \lambda_N^{N-1} \\ \lambda_1 & \lambda_2 & \cdots & \lambda_N \\ 1 & 1 & \cdots & 1 \end{pmatrix}$$

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- Conical defects have degenerate eigenvalues  $\rightarrow$  cannot be gauge-transformed into this form.

- If eigenvalues not degenerate, diagonalized by Vandermonde

$$V = \begin{pmatrix} \lambda_1^{N-1} & \lambda_2^{N-1} & \cdots & \lambda_N^{N-1} \\ \lambda_1 & \lambda_2 & \cdots & \lambda_N \\ 1 & 1 & \cdots & 1 \end{pmatrix}$$

- Conical defects have degenerate eigenvalues  $\rightarrow$  cannot be gauge-transformed into this form.
- Can falloff conditions be relaxed? E.g. include new sectors with falloff conditions

$$(A - A_{con.def.})|_{\rho \rightarrow \infty} = \mathcal{O}(1)$$

Other highest weight reps. of  $\mathcal{W}_N$ ?



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- what is higher spin geometry?